

A model for the present universe  
as a global attractor

Grigoris Panotopoulos, University of Crete

HEP2006, April 13-16, Ioannina

# Outline

- Introduction
- The model
- Analysis
- Conclusions

## Introduction

-According to observational data the Universe today is **accelerating**

- Within the framework of GR the second Friedmann equation:

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(1 + 3w)\rho$$

-Divide fluids into 2 categories:

a) Decelerating ( $w > -1/3$ ), b) Accelerating ( $w < -1/3$ )

- Simpler candidate: **Cosmological constant** (CC) (like a perfect fluid with  $w = -1 = \text{constant}$ )

- Other candidates: **Scalar fields** coupled to gravity (e.g. phantom, quintessence, with a varying  $w$ ) with some self-interaction potential

In general, dark energy models are divided into two broad categories:

-**Dynamical** dark energy, in which one modifies the matter side of the field equations

$$G_{\mu\nu} = 8\pi G(T_{\mu\nu} + T_{\mu\nu}^{dark})$$

-**Geometrical** dark energy, in which one modifies the geometric side of the field equations

$$(G_{\mu\nu} + G_{\mu\nu}^{dark}) = 8\pi GT_{\mu\nu}$$

- A simple modified-gravity model is the **DGP** brane-world model (as generalized to cosmology by Deffayet)
- We are confined on a 3d hypersurface, called the **brane**, embedded in a higher-dimensional spacetime and gravity resides in the whole **bulk**
- Brane-worlds are inspired from **String/M-theory** and contain the basic ingredients (extra dimensions, extended objects etc). They include corrections to usual Einstein-Hilbert term for gravity
- A IR correction is the **induced gravity** (IG) term on the brane (important at low energies)
- A UV correction is the **Gauss-Bonnet** (GB) term in the bulk (important at high energies, e.g. inflation)
- Interested in cosmic acceleration → IG correction

## The model

Action for the model

$$S = \int d^5x \sqrt{-g} (M^3 R - \Lambda) + \int d^4x \sqrt{-h} (m^2 \hat{R} - V)$$

plus matter in the bulk and perfect fluid on the brane

Einstein's equations

$$G_{AC} = \frac{1}{2M^3} T_{AC}|_{tot}$$

where

$$T_C^A|_{tot} = T_C^A|_{v,B} + T_C^A|_{m,B} + T_C^A|_{v,b} + T_C^A|_{m,b} + T_C^A|_{ind}$$

Seek for cosmological solutions

$$ds^2 = -n(t, y)^2 dt^2 + a(t, y)^2 \gamma_{ij} dx^i dx^j + b(t, y)^2 dy^2$$

Compute the **Einstein tensor**

$$\begin{aligned}
 G_{00} &= 3 \left\{ \frac{\dot{a}}{a} \left( \frac{\dot{a}}{a} + \frac{\dot{b}}{b} \right) - \frac{n^2}{b^2} \left[ \frac{a''}{a} + \frac{a'}{a} \left( \frac{a'}{a} - \frac{b'}{b} \right) \right] + \frac{kn^2}{a^2} \right\} \\
 G_{ij} &= \frac{a^2}{b^2} \gamma_{ij} \left\{ \frac{a'}{a} \left( \frac{a'}{a} + \frac{2n'}{n} \right) - \frac{b'}{b} \left( \frac{n'}{n} + \frac{2a'}{a} \right) + \frac{2a''}{a} + \frac{n''}{n} \right\} \\
 &\quad + \frac{a^2}{n^2} \gamma_{ij} \left\{ \frac{\dot{a}}{a} \left( \frac{2\dot{n}}{n} - \frac{\dot{a}}{a} \right) - \frac{2\ddot{a}}{a} + \frac{\dot{b}}{b} \left( \frac{\dot{n}}{n} - \frac{2\dot{a}}{a} \right) - \frac{\ddot{b}}{b} \right\} - k \gamma_{ij} \\
 G_{05} &= 3 \left( \frac{n' \dot{a}}{n a} + \frac{a' \dot{b}}{a b} - \frac{\dot{a}'}{a} \right) \\
 G_{55} &= 3 \left\{ \frac{a'}{a} \left( \frac{a'}{a} + \frac{n'}{n} \right) - \frac{b^2}{n^2} \left[ \frac{\ddot{a}}{a} + \frac{\dot{a}}{a} \left( \frac{\dot{a}}{a} - \frac{\dot{n}}{n} \right) \right] - \frac{kb^2}{a^2} \right\}
 \end{aligned}$$

Compute the various **energy-momentum tensors**

$$T_C^A|_{v,B} = \text{diag}(-\Lambda, -\Lambda, -\Lambda, -\Lambda, -\Lambda)$$

$$T_C^A|_{v,b} = \text{diag}(-V, -V, -V, -V, 0) \frac{\delta(y)}{b}$$

$$T_C^A|_{m,b} = \text{diag}(-\rho, p, p, p, 0) \frac{\delta(y)}{b}$$

$$T_0^0|_{ind} = \frac{6m^2}{n^2} \left( \frac{\dot{a}^2}{a^2} + \frac{kn^2}{a^2} \right) \frac{\delta(y)}{b}$$

$$T_j^i|_{ind} = \frac{2m^2}{n^2} \left( \frac{\dot{a}^2}{a^2} - \frac{2\dot{a}\dot{n}}{an} + \frac{2\ddot{a}}{a} + \frac{kn^2}{a^2} \right) \delta_j^i \frac{\delta(y)}{b}$$



Define new quantities

$$\begin{aligned}r_c &= \frac{m^2}{M^3} \\ \lambda &= \frac{2V}{m^2} + \frac{12}{r_c^2} - \frac{\Lambda}{M^3} \\ \mu &= \frac{V}{6m^2} + \frac{2}{r_c^2} \\ \gamma &= \frac{1}{12m^2} \\ \beta &= \frac{1}{\sqrt{3}r_c}\end{aligned}$$

Derive the cosmological equations on the brane

$$\dot{\rho} + 3H(1 + w)\rho = -T$$

$$H^2 = \mu + 2\gamma\rho + \beta\psi - \frac{k}{a^2}$$

$$\dot{\psi} + 2H\left(\psi - \frac{\lambda + 6(1-3w)\gamma\rho}{\psi}\right) = \frac{2\gamma T}{\beta}$$

$$\frac{\ddot{a}}{a} = \mu - (1+3w)\gamma\rho + \beta\frac{\lambda + 6(1-3w)\gamma\rho}{\psi}$$

where  $\psi \equiv \sqrt{\lambda + 24\gamma\rho + \chi}$  and assume  $T(\rho) = A\rho^\nu$ ,  $\nu > 0$

## Analysis

Define new variables (for flat universe,  $k = 0$ )

$$\omega_m = \frac{2\gamma\rho}{D^2} \quad , \quad \omega_\psi = \frac{\beta\psi}{D^2} \quad , \quad Z = \frac{H}{D}$$

New equations

$$\omega_m + \omega_\psi = 1$$

$$\begin{aligned} \omega'_m = \omega_m & \left[ (1+3w)(\omega_m-1)Z - \frac{A}{\sqrt{|\mu|}} \left( \frac{|\mu|\omega_m}{2\gamma} \right)^{\nu-1} (1-Z^2)^{\frac{3}{2}-\nu} \right. \\ & \left. - 2Z(1-Z^2) \frac{1-Z^2-3(1-3w)\beta^2\mu^{-1}\omega_m}{1-\omega_m} \right] \\ Z' = (1-Z^2) & \left[ (1-Z^2) \frac{1-Z^2-3(1-3w)\beta^2\mu^{-1}\omega_m}{1-\omega_m} - 1 \right. \\ & \left. - \frac{1+3w}{2}\omega_m \right] \end{aligned}$$

## 1. Fixed points:

- Fixed point: If initial condition, the system stays there. It is found by setting derivatives equal to 0
  - Character: **Attractor, repeller, saddle** (depends on the characteristic matrix)
  - **Attractor**: 2 eigenvalues (real part) negative
  - Repeller**: 2 eigenvalues (real part) positive
  - Saddle**: 1 eigenvalue positive, 1 negative
- The results are summarized in the tables

	$\nu < 3/2$	$\nu = 3/2$	$\nu > 3/2$
No. of F.P.	1	0 or 1	1
Nature	A	A	S

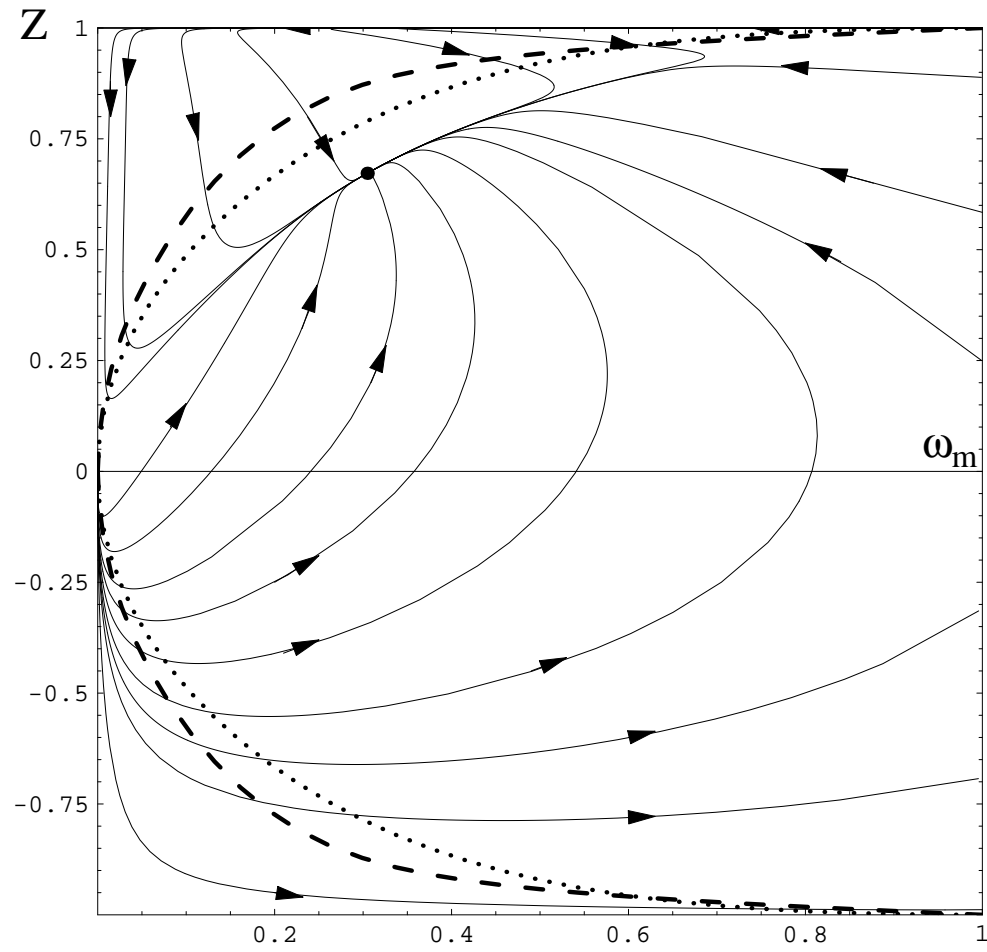
Table 1: The fixed points for  $w=0$ , influx

	$\nu < 3/2$	$\nu = 3/2$	$3/2 < \nu < 2$	$\nu = 2$	$\nu > 2$
No. of F.P.	1	0 or 1	0 or 2	0 or 1	1
Nature	A	A	A,S	S	S

Table 2: The fixed points for  $w=1/3$ , influx

## 2. Phase-portraits ( $\omega_m - Z$ plane)

The history of the universe can be shown



Influx,  $w=0$ ,  $\nu < 3/2$

3. **Dark energy:** It is a result of the geometry and the energy exchange

We can write a semi-conservation law for dark energy

$$\dot{\rho}_{DE} + 3(1+w_{DE})H\rho_{DE} = T$$

where  $w_{DE}$  is

$$w_{DE} = \frac{-1}{3(1-\omega_m)} \left[ 2Z^2 - \omega_m - 1 - 6(1-3w) \frac{\beta^2 \omega_m (1-Z^2)}{\mu (Z^2 - \omega_m)} \right]$$

**Evolving** in time

At the **fixed point** we obtain

$$w_{DE*} = -1 - \frac{1+w}{\Omega_{m*}^{-1} - 1}$$

and for  $\Omega_{m*} = 0.3$ ,  $w_{DE*} = -1.4$



## Conclusions

- In our scenario, the present universe can be easily realized as a unique attractor of the dynamical system
- The dark energy is a result of the geometry and the energy exchange
- The parameter  $w$  of the dark energy evolves in time and crosses the  $w = -1$  barrier from larger to lower values
- For  $\Omega_M = 0.3$ , the parameter  $w$  of the dark energy is today  $w_0 = -1.4$ , independently of the parameters of the models